

Reading assignment: pages 173-184 and 197-203 of Fetter and Walecka.

Problem 1. A relativistic particle with mass m and position x moving in a static potential $V(x)$ in one dimension is described by the Lagrangian $L(x, \dot{x}) = -mc^2 \sqrt{1 - \dot{x}^2/c^2} - V(x)$, where c is the speed of light.

- (a) Find the Euler-Lagrange equation for $x(t)$.
- (b) Find the canonical momentum p and write a simplified expression for the Hamiltonian $H(x, p)$. Is H a constant of the motion?

Problem 2. A particle with mass m , position \vec{r} , and momentum \vec{p} has angular momentum $\vec{L} = \vec{r} \times \vec{p}$. It is useful to define the antisymmetric symbol ϵ_{ijk} as

$$\epsilon_{123} = \epsilon_{231} = \epsilon_{312} = 1, \quad \epsilon_{132} = \epsilon_{213} = \epsilon_{321} = -1, \quad (1)$$

and all other components $\epsilon_{ijk} = 0$. It obeys the identities

$$\begin{aligned} \epsilon_{ijk}\epsilon_{lnk} &= \delta_{il}\delta_{jn} - \delta_{in}\delta_{lj}, \\ \epsilon_{ijk}\epsilon_{ljk} &= 2\delta_{il} \\ \epsilon_{ijk}\epsilon_{ijk} &= 6. \end{aligned}$$

Now the components of the angular momentum can be written as

$$L_i = \epsilon_{ijk}r_jp_k.$$

Note that the Einstein summation convention is in effect: repeated indices are summed over except when they appear on both sides of an equation. Evaluate the following Poisson brackets in fully simplified form:

- (a) $[r_j, L_k]_{\text{PB}} = ?$
- (b) $[p_j, L_k]_{\text{PB}} = ?$
- (c) $[L_j, L_k]_{\text{PB}} = ?$
- (d) $[r_j, L^2]_{\text{PB}} = ?$

Problem 3. A particle of mass m and charge q moves in a background magnetic field \vec{B} .

- (a) Show that $[m\dot{r}_i, m\dot{r}_j]_{\text{PB}} = q\epsilon_{ijk}B_k$. Note that this means that the Poisson brackets of two *kinematic momenta* (also known as mechanical momenta) can be non-zero in the presence of a magnetic field orthogonal to them both, even though the Poisson bracket of any two *canonical momenta* is always zero.
- (b) Compute $[m\dot{r}_i, r_j]_{\text{PB}} = ?$

Problem 4. Consider a system with phase space coordinates q_1, q_2 and their canonical conjugate momenta p_1, p_2 , with Hamiltonian $H = \frac{1}{2}(q_1^2 + q_2^2 + p_1^2 + p_2^2)$. Consider the new coordinates Q_1, Q_2 and P_1, P_2 on phase space, defined by

$$\begin{aligned} q_1 &= Q_1 \cos \alpha + P_2 \sin \alpha, \\ q_2 &= Q_2 \cos \beta + P_1 \sin \beta, \\ p_1 &= -Q_2 \sin \beta + P_1 \cos \beta, \\ p_2 &= -Q_1 \sin \alpha + P_2 \cos \alpha, \end{aligned}$$

where α and β are constants.

- (a) Find the inverse relations, solving for Q_1, Q_2, P_1, P_2 in terms of q_1, q_2, p_1, p_2 . Make sure your answers are in fully simplified form.
- (b) Compute all of the Poisson brackets of every pair from Q_1, Q_2, P_1, P_2 . By requiring that the relation is a canonical transformation, solve for β in terms of α .
- (c) Express each of p_1, p_2, P_1 , and P_2 in terms of q_1, q_2, Q_1, Q_2 . Use these results to find the Type 1 generating function $F(q_1, q_2, Q_1, Q_2)$ for the canonical transformation. [For the definition of this type of generating function, see Fetter and Walecka, eqs. (34.9)-(34.11).]
- (d) Express each of p_1, p_2, Q_1 , and Q_2 in terms of q_1, q_2, P_1, P_2 . Use these results to find the Type 2 generating function $F_2(q_1, q_2, P_1, P_2)$ for the canonical transformation, which means that

$$p_n = \frac{\partial F_2}{\partial q_n}, \quad Q_n = \frac{\partial F_2}{\partial P_n}.$$

- (e) Find the new Hamiltonian $H(Q_1, Q_2, P_1, P_2)$ in fully simplified form. What are the Hamilton equations of motion for Q_1, Q_2, P_1, P_2 ?
- (f) Use your results to find the solutions for q_1, q_2, p_1, p_2 as functions of the time t , under the constraints that $Q_2 = 0$ and $P_2 = 0$ for all times, and with the boundary conditions $q_1(0) = x_0$ and $\dot{q}_1(0) = v_0$. [You do not need to introduce a Lagrange multiplier for the constraints. Your answers should depend only on x_0, v_0, α , and t .]

Problem 5. A group of particles, all with the same mass m , have initial heights z and initial vertical momenta p lying in the rectangle $-a < z < a$ and $-b < p < b$. The particles then fall freely for a time t in a uniform gravitational field with acceleration g downwards. Find the region in phase space that they occupy at time t , sketch it, and show by direct calculation (not appealing to any fancy general theorems) that the phase space area is still $4ab$. This is an illustration of Liouville's Theorem, which says that areas in phase space are preserved under Hamiltonian time evolution.