

Quiz #2

Solutions

① The radial probability density for the 1s state is

$$P(r) = r^2 |R_{10}(r)|^2 = \frac{4}{a_0^3} r^2 e^{-2r/a_0}$$

(a) The maximum of $P(r)$ occurs when $\frac{dP}{dr} = 0$

$$\begin{aligned} \frac{dP}{dr} = 0 &= \left(\frac{4}{a_0^3}\right) \left[2r e^{-2r/a_0} - \frac{2r^2}{a_0} e^{-2r/a_0} \right] \\ &= \left(\frac{4}{a_0^3}\right) 2r e^{-2r/a_0} \left[1 - \frac{r}{a_0} \right] \end{aligned}$$

$$\Rightarrow \left(1 - \frac{r}{a_0}\right) = 0 \quad \Rightarrow r = a_0 \quad \text{at the maximum}$$

$$(b) P(r) = r^2 |R_{20}|^2 = \frac{1}{(2a_0)^3} r^2 \left(2 - \frac{r}{a_0}\right)^2 e^{-r/a_0}$$

$$\frac{dP}{dr} = 0 = \frac{1}{(2a_0)^3} \left[2r \left(2 - \frac{r}{a_0}\right)^2 e^{-r/a_0} + 2r^2 \left(2 - \frac{r}{a_0}\right) \left(-\frac{1}{a_0}\right) e^{-r/a_0} - \left(\frac{1}{a_0}\right) r^2 \left(2 - \frac{r}{a_0}\right)^2 e^{-r/a_0} \right]$$

$$= \frac{1}{(2a_0)^3} r \left(2 - \frac{r}{a_0}\right) e^{-r/a_0} \left[2 \left(2 - \frac{r}{a_0}\right) + 2r \left(-\frac{1}{a_0}\right) - \left(\frac{1}{a_0}\right) r \left(2 - \frac{r}{a_0}\right) \right]$$

$$= \frac{1}{(2a_0)^3} r \left(2 - \frac{r}{a_0}\right) \left[4 - \frac{2r}{a_0} - \frac{2r}{a_0} - \frac{2r}{a_0} + \frac{r^2}{a_0^2} \right]$$

$$0 = \frac{1}{(2a_0)^3} r \left(2 - \frac{r}{a_0}\right) \left[\frac{r^2}{a_0^2} - \frac{6r}{a_0} + 4 \right] = \frac{1}{(2a_0)^3} \frac{1}{a_0^2} r \left(2 - \frac{r}{a_0}\right) [r^2 - 6a_0r + 4a_0^2]$$

The maximum cannot occur when $r(2 - r/a_0) = 0$ since at these points ~~the~~ $R_{20} \approx r(2 - r/a_0)e^{-r/a_0} = 0$. Thus, the maximum must occur when

$$r^2 - 6a_0r + 4a_0^2 = 0$$

$(r=0, r=2a_0)$ give minima

$$\Rightarrow r = \frac{6a_0 \pm \sqrt{(6a_0)^2 - 4(4a_0^2)}}{2}$$

$$= \frac{3a_0 \pm \sqrt{20a_0^2}}{2} = 3a_0 \pm a_0\sqrt{5}$$

$$= (3 \pm \sqrt{5})a_0 = (5.236 a_0 \text{ and } 1.586 a_0)$$

two maxima

②

Recall that the expectation value of a function $f(x)$ in one-dimensions is

$$\langle f(x) \rangle = \int_{-\infty}^{\infty} f(x) |\psi(x)|^2 dx$$

In 3-dimensions, we integrate over all space

$$\langle f(r, \theta, \phi) \rangle = \int_{\text{all space}} f(r, \theta, \phi) |\psi(r, \theta, \phi)|^2 dV$$

Thus, the expectation value of $\frac{1}{r}$ is (for the 1s state)

$$\begin{aligned} \text{(a)} \quad \langle \frac{1}{r} \rangle &= \int_{\text{all space}} \left(\frac{1}{r} \right) \left| \frac{1}{\sqrt{\pi a_0^3}} e^{-r/a_0} \right|^2 dV \\ &= \int_{r=0}^{\infty} \left(\frac{1}{r} \right) \left(\frac{1}{\pi a_0^3} \right) e^{-2r/a_0} r^2 dr \int_{\theta=0}^{\pi} \sin \theta d\theta \int_{\phi=0}^{2\pi} d\phi \end{aligned}$$

Note that: $\int_0^{2\pi} d\phi = \phi \Big|_0^{2\pi} = 2\pi - 0 = 2\pi$

$$\int_0^{\pi} \sin \theta d\theta = -\cos \theta \Big|_0^{\pi} = -(\cos \pi - \cos 0) = -(-1 - 1) = 2$$

The integral over angles then gives a factor $(2)(2\pi) = 4\pi$
and

$$\begin{aligned} \langle \frac{1}{r} \rangle &= \frac{(4\pi)}{\pi a_0^3} \int_0^{\infty} r e^{-2r/a_0} dr \\ &= \left(\frac{4}{a_0^3} \right) \left[\frac{1!}{(2/a_0)^2} \right] = \left(\frac{4}{a_0^3} \right) \left(\frac{a_0^2}{4} \right) = \frac{1}{a_0} \end{aligned}$$

(b) Since $U = -\frac{1}{4\pi\epsilon_0} \frac{e^2}{r}$

$$\text{then } \langle U \rangle = \left\langle -\frac{1}{4\pi\epsilon_0} \frac{e^2}{r} \right\rangle = -\frac{e^2}{4\pi\epsilon_0} \langle \frac{1}{r} \rangle = -\frac{1}{4\pi\epsilon_0} \frac{e^2}{a_0}$$

$$\text{since } \langle \frac{1}{r} \rangle = \frac{1}{a_0}$$

(c) Given that $\langle E \rangle = \langle K \rangle + \langle U \rangle$

Then $\langle K \rangle = \langle E \rangle - \langle U \rangle$

For the 1s state: $E = E_1 = -\frac{\hbar^2}{2ma_0^2}$

The expectation value of a constant is simply the constant since Ψ_{100} is normalized:

$$\begin{aligned}\langle E \rangle &= \int_{\text{all space}} E_1 |\Psi_{100}|^2 dV \\ &= E_1 \int |\Psi_{100}|^2 dV \quad (\text{since } E_1 = \text{constant}) \\ &= E_1 \quad (\text{since } \Psi_{100} \text{ is normalized})\end{aligned}$$

Then

$$\begin{aligned}\langle K \rangle &= -\frac{\hbar^2}{2ma_0^2} - \left(-\frac{1}{4\pi\epsilon_0} \frac{e^2}{a_0} \right) \\ &= \frac{-\hbar^2}{2ma_0 \left(4\pi\epsilon_0 \frac{\hbar^2}{me^2} \right)} + \frac{1}{4\pi\epsilon_0} \frac{e^2}{a_0} \quad \left(\text{since } a_0 = 4\pi\epsilon_0 \frac{\hbar^2}{me^2} \right) \\ &= -\frac{1}{4\pi\epsilon_0} \frac{e^2}{2a_0} + \frac{1}{4\pi\epsilon_0} \frac{e^2}{a_0} \\ &= \frac{1}{4\pi\epsilon_0} \frac{e^2}{2a_0} = \frac{\hbar^2}{2ma_0^2}\end{aligned}$$

3 (a) Schrodinger Equation is

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + \frac{1}{2} Cx^2 \psi = E\psi$$

For $\psi(x) = Ae^{-bx^2}$

$$\frac{d\psi}{dx} = -2bx \cdot Ae^{-bx^2}$$

$$\frac{d^2\psi}{dx^2} = -2bAe^{-bx^2} + 4b^2x^2Ae^{-bx^2}$$

Plugging into the Schrodinger Eq. gives

~~$$\frac{\hbar^2}{2m} [-2b + 4b^2x^2] Ae^{-bx^2} + \frac{1}{2} Cx^2 (Ae^{-bx^2}) = E (Ae^{-bx^2})$$~~

$$\left(\frac{1}{2} C - \frac{4b^2\hbar^2}{2m} \right) x^2 + \left(\frac{2b\hbar^2}{2m} - E \right) = 0$$

$$Bx^2 + D = 0$$

For $D=0$: $\frac{2b\hbar^2}{2m} - E = 0 \Rightarrow E = \frac{b}{m} \hbar^2$

For $B=0$: $\frac{1}{2} C - \frac{4b^2\hbar^2}{2m} = 0 \Rightarrow b = \frac{\sqrt{Cm}}{2\hbar}$

Then: $E = \frac{\hbar^2}{m} \left(\frac{\sqrt{Cm}}{2\hbar} \right) = \frac{1}{2} \hbar \sqrt{\frac{C}{m}} = \frac{1}{2} \hbar \omega_0$

Where $\omega_0 =$ resonance frequency

(b)

