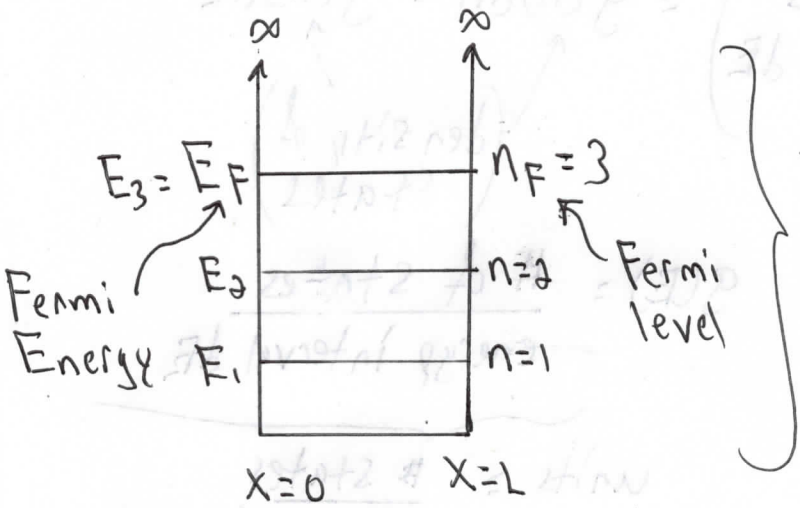


# Density & States

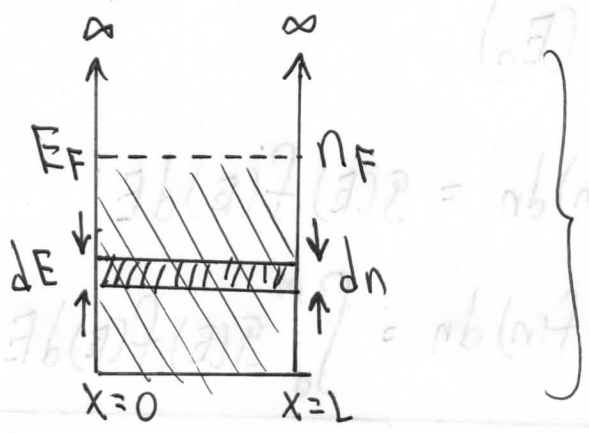
## Particle in a Box (infinite square well)



### Discrete Case

There are  
 $N = 3 \times 2 = 6$  states  
 ↑ ↑  
 energy 2 spins  
 levels each

Note:  $E_n = \left( \frac{\hbar^2 \pi^2}{2mL^2} \right) n^2 \Rightarrow$  as  $L \rightarrow \infty$ ,  $\Delta E_n \rightarrow 0$   
 density of states increases as  $L$  increases



Continuous case  
 for Solids, we have  
 Avogadro's # of energy levels  
 $N_A \sim 10^{23}$  levels

How many states are in the interval  $dn$  or  $dE$ ?  
 To answer this, let's define a density of States

Discrete: Number of states =  $N = \sum_{n=1}^3 dn = 2 + 2 + 2 = 6$

Krane's notation  $\rightarrow$  (degeneracy of each state) example above

Continuous:

$$dN = \left( \begin{array}{l} \text{Number of States} \\ \text{in interval } dn \text{ or } dE \end{array} \right) = g(n)dn = g(E)dE$$

(density of states)

$$g(n) = \frac{\# \text{ of states}}{\text{unit interval } dn}$$

units = # of states

$$g(E) = \frac{\# \text{ of states}}{\text{energy interval } dE}$$

units =  $\frac{\# \text{ states}}{\text{Joules (Energy)}}$

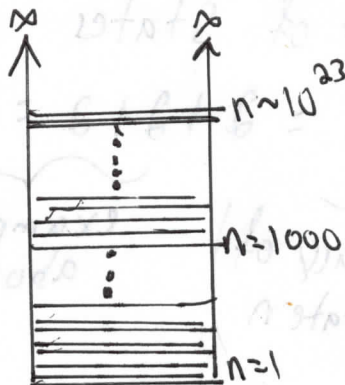
There is a probability  $f(E)$  that a state may be occupied

Discrete:  $N = \sum_n dn f(E_n)$

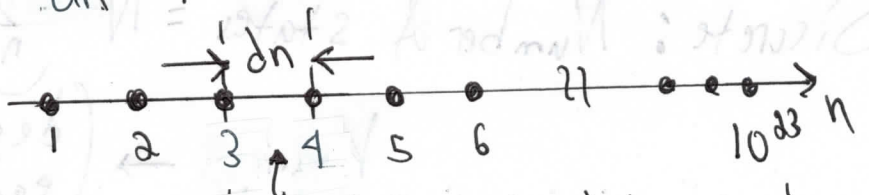
Continuous:  $dN = g(n) f(n) dn = g(E) f(E) dE$

$$N = \int_0^\infty g(n) f(n) dn = \int_0^\infty g(E) f(E) dE$$

1-Dimensional Example: Infinite square well in 1-dimension



We can plot all  $n$  states along an  $n$ -axis



notice, there is only 1 state in interval  $dn$

What is the density of states  $g(E)$ ? (3)

$$g(E)dE = g(n)dn$$

$$g(n) = \frac{1 \text{ state}}{\text{interval } dn} \times 2 = 2$$

↑ degeneracy factor  
for spin up & down states

Then,  $g(E) = g(n) \frac{dn}{dE} = 2 \frac{dn}{dE}$

For an infinite square well,  $E_n = \left( \frac{\hbar^2 \pi^2}{2mL^2} \right) n^2 = E_1 n^2$

Then,  $dE = d(E_1 n^2) = 2E_1 n dn$

$$\Rightarrow \frac{dn}{dE} = \frac{1}{2E_1 n}$$

$$E = E_1 n^2 \Rightarrow n = \sqrt{\frac{E}{E_1}}$$

Then,  $\frac{dn}{dE} = \frac{1}{2E_1 n} = \frac{1}{2E_1} \sqrt{\frac{E_1}{E}} = \frac{1}{2\sqrt{E_1 E}}$

This gives:  $g(E) = 2 \frac{dn}{dE} = \frac{1}{\sqrt{E_1 E}} = \left( \frac{L\sqrt{2m}}{\hbar\pi} \right) \frac{1}{\sqrt{E}}$

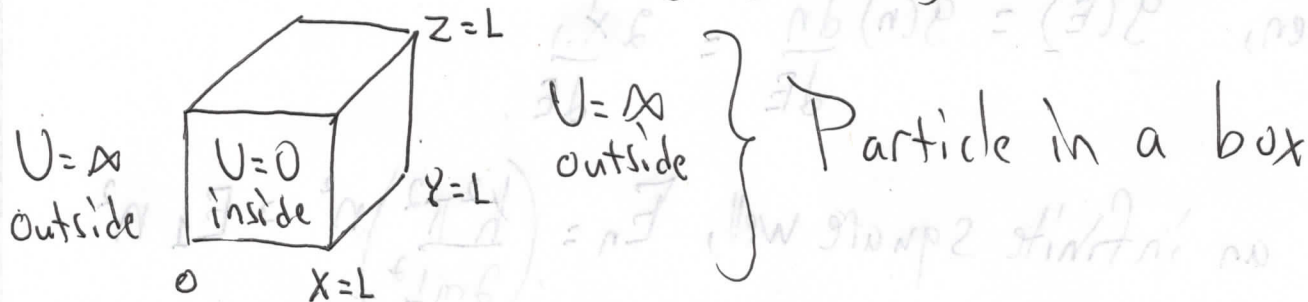
(# of states in interval  $dE$ ) =  $dN = g(E) f(E) dE$   
=  $\left( \frac{L\sqrt{2m}}{\hbar\pi} \right) \frac{1}{\sqrt{E}} f(E) dE$  } 1-dimension



Since  $g(E) \rightarrow \infty$  as  $L \rightarrow \infty$ , textbooks like to define a density of states per unit length (4)

$$g_L(E) = \frac{g(E)}{L} = \left( \frac{\sqrt{2m}}{h\pi} \right) \frac{1}{\sqrt{E}} f(E) dE$$

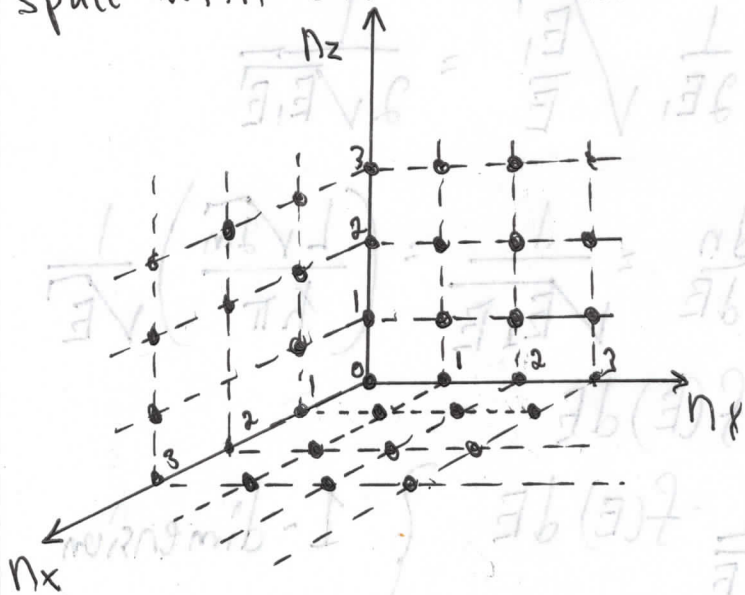
3-dimensional Example: Particle in a Box  
 $U=0$  inside,  $U=\infty$  outside box



The allowed energies are similar to that for a 1-dimensional problem:

$$E_{n_x n_y n_z} = \left( \frac{h^2 \pi^2}{2mL^2} \right) (n_x^2 + n_y^2 + n_z^2)$$

Each state can be defined as a point in 3-dimensional space with axes  $n_x, n_y, n_z$

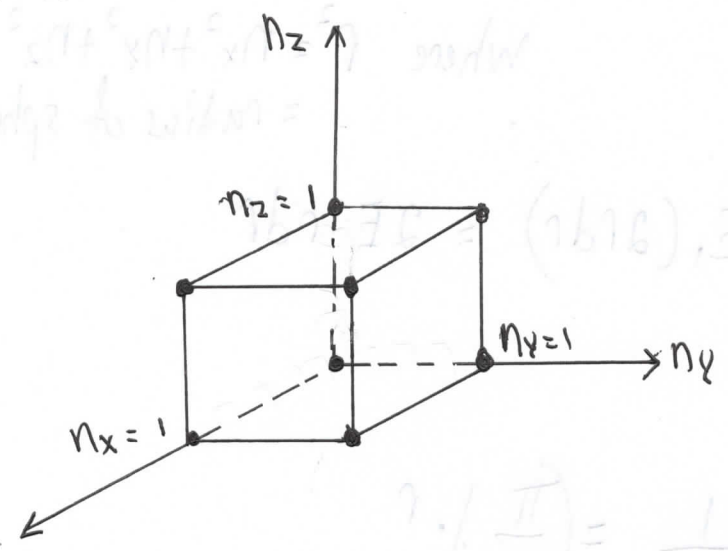


only positive values of  $n_x, n_y, n_z$  are allowed:

$$n_x \geq 0, n_y \geq 0, n_z \geq 0$$

Thus, only the 1<sup>st</sup> quadrant in space is filled with allowed states.

Let's examine the states where  $n_x = 1, 0$ ,  $n_y = 1, 0$ ,  $n_z = 1, 0$



Note that there is only 1 state in this cube (the corners are shared by 8 other cubes)

spin up & down states

Thus,  $g(n) = \frac{\# \text{ of states}}{\text{unit volume } dn_x dn_y dn_z} = 1 \text{ state} \times 2 = 2$

$g(E)dE = g(n) dV_n = 2 \cdot d n_x d n_y d n_z$

Volume in n-space

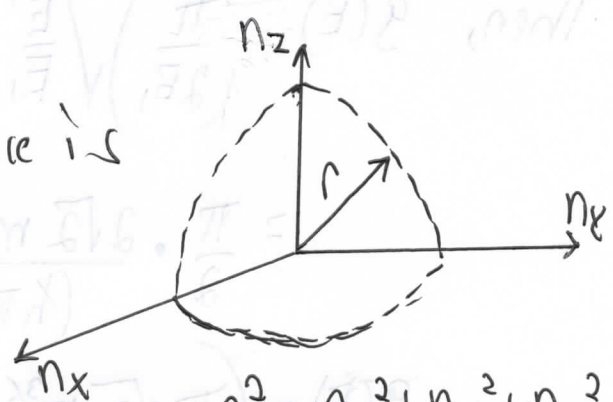
A spherical volume in n-space is

$$V_n = \frac{4}{3} \pi r^3$$

$$dV_n = d\left(\frac{4}{3} \pi r^3\right)$$

$$= \frac{4}{3} \pi (3r^2 dr)$$

$$dV_n = 4\pi r^2 dr$$



$r^2 = n_x^2 + n_y^2 + n_z^2$

$r = \text{radius of a sphere}$

$\Rightarrow g(E)dE = 2 \cdot (4\pi r^2 dr) \cdot \frac{1}{8}$

$g(E)dE = \pi r^2 dr$

$\Rightarrow g(E) = \pi r^2 \frac{dr}{dE}$

$n_x, n_y, n_z$  only allowed in 1<sup>st</sup> quadrant of sphere =  $\frac{1}{8}$ th of the volume of the sphere

$n_x, n_y, n_z$  must not be negative

Since  $E = \underbrace{\left(\frac{\hbar^2 \pi^2}{2mL^2}\right)}_{E_1} (n_x^2 + n_y^2 + n_z^2) = E_1 r^2$

where  $r^2 = n_x^2 + n_y^2 + n_z^2$   
= radius of sphere

$$dE = d(E_1 r^2) = E_1 (2r dr) = 2E_1 r dr$$

$$\Rightarrow \frac{dr}{dE} = \frac{1}{2E_1 r}$$

$$\Rightarrow g(E) = \pi r^2 \frac{dr}{dE} = \pi r^2 \cdot \frac{1}{2E_1 r} = \left(\frac{\pi}{2E_1}\right) \cdot r$$

But,  $E = E_1 r^2 \Rightarrow r = \sqrt{\frac{E}{E_1}}$

Then,  $g(E) = \left(\frac{\pi}{2E_1}\right) \sqrt{\frac{E}{E_1}} = \left(\frac{\pi}{2E_1^{3/2}}\right) \sqrt{E} = \frac{\pi}{2} \left(\frac{2mL^2}{\hbar^2 \pi^2}\right)^{3/2} \sqrt{E}$

$$= \frac{\pi}{2} \cdot \frac{2\sqrt{2} m^{3/2} L^3}{(\hbar \pi)^3} \cdot \sqrt{E}$$

$$g(E) = \left(\frac{\pi \sqrt{2} m^{3/2} V}{(\hbar \pi)^3}\right) \sqrt{E}$$

where  $V = L^3 = \text{Volume}$   
in  $x, y, z$ -space

(# of states in interval  $dE$ ) =  $dN = g(E) f(E) dE$   
=  $\left(\frac{\pi \sqrt{2} m^{3/2} V}{(\hbar \pi)^3}\right) \sqrt{E} f(E) dE$  } 3-dimensions

(density of states per unit volume) =  $g_V(E) = \left(\frac{\pi \sqrt{2} m^{3/2}}{(\hbar \pi)^3}\right) \sqrt{E}$  ← Krane's notation